

Interpretations of the wave function

We are familiar with the probability density

$$\rho(\vec{x}, t) = |\Psi(\vec{x}, t)|^2 = |\langle \vec{x} | \alpha, t_0; t \rangle|^2$$

↓
Introduce the probability flux \Rightarrow

$$\vec{j}(\vec{x}, t) = -\frac{i\hbar}{2m} [\Psi^* \vec{\nabla} \Psi - (\vec{\nabla} \Psi^*) \Psi] =$$

$$\stackrel{\uparrow}{=} -\frac{i\hbar}{2m} \cdot 2i \operatorname{Im}(\Psi^* \vec{\nabla} \Psi) = \frac{\hbar}{m} \operatorname{Im}(\Psi^* \vec{\nabla} \Psi)$$

$$a+ib - (a-ib) = 2ib$$

Is there a connection between ρ and \vec{j} ? \Rightarrow

$$\frac{\partial \rho}{\partial t} = \frac{\partial}{\partial t} (\Psi^* \Psi) = \underbrace{\frac{\partial \Psi^*}{\partial t} \Psi}_{\frac{i}{\hbar} H \Psi^*} + \Psi^* \underbrace{\frac{\partial \Psi}{\partial t}}_{-\frac{i}{\hbar} H \Psi} = \frac{i}{\hbar} (H \Psi^*) \Psi - \Psi^* (H \Psi)$$

$$\ominus \frac{i}{\hbar} \left(-\frac{\hbar^2}{2m} (\Delta \Psi^*) \Psi + \cancel{(\vec{\nabla} \Psi^*) \Psi} + \Psi^* \frac{\hbar^2}{2m} (\Delta \Psi) - \cancel{\Psi^* (\vec{\nabla} \Psi)} \right)$$

$$= \frac{i}{\hbar} \cdot \frac{\hbar^2}{2m} (\psi^* \Delta \psi - \psi \Delta \psi^*) = \frac{i\hbar}{2m} \cdot 2i \operatorname{Im}(\psi^* \Delta \psi) \quad (2)$$

$$= -\frac{\hbar}{m} \operatorname{Im}(\psi^* \Delta \psi)$$

Go back to $\vec{j}(\vec{x}, t)$ and consider $\vec{\nabla} \cdot \vec{j} \Rightarrow$

$$\vec{\nabla} \cdot \vec{j} = \frac{\hbar}{m} \vec{\nabla} [\operatorname{Im}(\psi^* \vec{\nabla} \psi)] = -\frac{i\hbar}{2m} \vec{\nabla} [\psi^* \vec{\nabla} \psi - (\vec{\nabla} \psi^*) \psi]$$

$$= -\frac{i\hbar}{2m} [\cancel{\vec{\nabla} \psi^* \cdot \vec{\nabla} \psi} + \psi^* \Delta \psi - (\Delta \psi^*) \psi - \cancel{(\vec{\nabla} \psi^*)(\vec{\nabla} \psi)}] = \frac{\hbar}{m} \operatorname{Im}[(\Delta \psi) \psi^*]$$

$$\boxed{\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot \vec{j} = 0}$$

\Rightarrow continuity equation

\Rightarrow conservation of the probability

\Downarrow compare with charge conserv. in E3M, fluid dynamics etc.

What is the physical meaning of \vec{j} ?

\Downarrow
We know that $\int_{V_0} d\vec{x} \rho(\vec{x}, t)$ is a probability to find a particle in volume V_0 ,

$$\int_{\text{all space}} \rho(\vec{x}, t) d\vec{x} = 1$$

What about $\int \vec{j}(\vec{x}, t) d\vec{x} = -\frac{i\hbar}{2m} \int_{\text{all space}} [\psi^* \vec{\nabla} \psi - (\vec{\nabla} \psi^*) \psi] d\vec{x} =$ (3)

$= \frac{1}{2m} \int_{\text{all space}} [\psi^* (-i\hbar \vec{\nabla} \psi) + (i\hbar \vec{\nabla} \psi^*) \psi] d\vec{x} =$

$\underbrace{\hspace{10em}}_{\vec{P}\psi} \quad \quad \quad \underbrace{\hspace{10em}}_{\vec{P}^*\psi^*}$

$= \frac{1}{2m} \int_{\text{all space}} (\psi^* (\vec{P}\psi) + (\vec{P}^*\psi^*) \psi) d\vec{x} =$

\downarrow
 $a+ib + a-ib = 2a$

$= \frac{1}{2m} \cdot 2 \operatorname{Re} \int_{\text{all space}} \psi^*(\vec{x}, t) \vec{P} \psi(\vec{x}, t) d\vec{x} =$

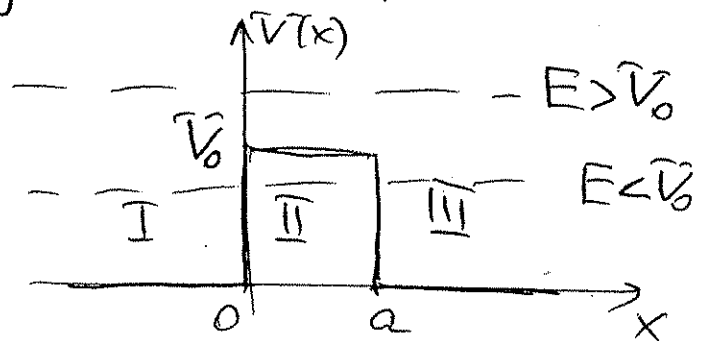
$= \frac{1}{m} \langle \vec{P} \rangle_t$ ← expectation value of the momentum operator at time t

Tunneling

Consider a particle moving in $V(x)$:

$$V(x) = \begin{cases} 0, & x < 0 \\ V_0, & 0 < x < a \\ 0, & x > a \end{cases}$$

$(V_0 > 0)$



Classically: particle reflected if $E < V_0$
and transmitted if $E > V_0$

QM: ? \Rightarrow 1D scattering problem

$$\Psi_{\text{I}}(x) = A e^{ikx} + B e^{-ikx} \quad (x < 0)$$

$$\Psi_{\text{III}}(x) = C e^{ikx} + D e^{-ikx} \quad (x > a)$$

$$k = \sqrt{\frac{2mE}{\hbar^2}}$$

$\Psi_{\text{II}}(x) \rightarrow$ depends on whether $E > V_0$ or $E < V_0$

Let's specify that the particle is incident on the barrier from the left \Rightarrow then there is nothing at large $x > 0$ to cause a reflection (i.e. e^{-ikx} term at $x > a$) $\Rightarrow D = 0$

Now look for transmission and reflection characteristics of the barrier:

reflection coefficient $R = \left| \frac{B}{A} \right|^2$

transmission coefficient $T = \left| \frac{C}{A} \right|^2$

\Rightarrow

Probability current densities:

$$x < 0: \quad j = -\frac{i\hbar}{2m} \left[\psi^* \frac{d\psi}{dx} - \psi \frac{d\psi^*}{dx} \right] =$$

$$= \frac{\hbar}{m} \operatorname{Im} \left[\psi^* \frac{d\psi}{dx} \right] = \frac{\hbar}{m} k (|A|^2 - |B|^2)$$

$$(A e^{-ikx} + B e^{ikx}) \cdot ik (A e^{ikx} - B e^{-ikx}) =$$

$$= ik (|A|^2 - |B|^2 - A^* B e^{-2ikx} + B^* A e^{2ikx})$$

$$x > a: \quad j = \frac{\hbar k}{m} |C|^2$$

$$\begin{aligned} a + ib - a + ib &= 2ib = \\ &= 2i \operatorname{Im} [B^* A e^{2ikx}] \end{aligned}$$

$$v = \frac{\hbar k}{m}$$

↑
particle velocity

$$v |A|^2$$

↑
incident intensity

$$v |B|^2$$

↑
reflected intensity

$$v |C|^2$$

↑
transmitted intensity

reflection coefficient $R = \frac{|B|^2}{|A|^2}$

transmission coefficient $T = \frac{|C|^2}{|A|^2}$

Independent of the normalization of ψ

1) $E < V_0$

$$\psi_{II}(x) = F e^{\alpha x} + G e^{-\alpha x}, \quad 0 < x < a$$

$$\alpha = \sqrt{\frac{2m(V_0 - E)}{\hbar^2}}$$

Match boundary conditions:

$\psi(x)$, $\psi'(x)$ are continuous at $x=0$ and $x=a \Rightarrow$

$$R = \left[1 + \frac{4E(V_0 - E)}{V_0^2 \sinh^2 \kappa a} \right]^{-1}$$

$$T = \left[1 + \frac{V_0^2 \sinh^2 \kappa a}{4E(V_0 - E)} \right]^{-1}$$

Show this for
Homework!

$R + T = 1$ (conservation of the probability flux)

$T \neq 0 \Rightarrow$ barrier penetration, or tunneling

If $E \rightarrow 0 \Rightarrow T \rightarrow 0$

If $E \rightarrow V_0$ (but $E < V_0$) $\Rightarrow \kappa a = \sqrt{\frac{2m(V_0 - E)}{\hbar^2}} a =$

$$\Rightarrow \sinh^2 \kappa a \approx (\kappa a)^2 = \frac{2m}{\hbar^2} (V_0 - E) a^2$$

$$\text{Then } T = \left[1 + \frac{V_0 m a^2}{2\hbar^2} \right]^{-1}$$

\uparrow a measure of the
"opacity" of the barrier
(the larger V_0 and the
wider $a \Rightarrow$ smaller T)

Classical limit:

$$\hbar \rightarrow 0 \Rightarrow T = 0!$$

2) $E > V_0$

$$\Psi_{II}(x) = Fe^{ik'x} + Ge^{-ik'x}, \quad 0 < x < a$$

$$k' = \frac{\sqrt{2m(E-V_0)}}{\hbar}$$

$$R = \left[1 + \frac{4E(E-V_0)}{V_0^2 \sin^2 k'a} \right]^{-1}$$

$$T = \left[1 + \frac{V_0^2 \sin^2 k'a}{4E(E-V_0)} \right]^{-1}$$

(derive (homework)!) }

Important: $T < 1$! \leftarrow contradiction ⁱⁿ to the classical prediction!!

$T = 1$ only if $k'a = \pi + \pi n = \pi n$
 \uparrow integer
 condition of destructive interference between the reflections at $x=0$ and $x=a$

At $E \gg V_0 \Rightarrow T \rightarrow 1$

- Physical examples of tunneling:
- α -particle emission from a nucleus
 - electron emission from metal
 - scanning tunneling microscope
- See Nature paper \Rightarrow

