Feynman's Quantum Paths

(Advanced ⇒ Relativity, Quantum Chromodynamics)

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Course: Computational Physics II



Generalize Classical Trajectory to QM Probability

- Classical Mech: single x(t) path $=\bar{x}$
- QM: waves = statistical, no path
- Dirac: Hamilton's least-action prin

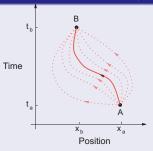


- F: Look for quantum least-action principle
- Hamilton: space-time path variation δ calculus
- F: quantum particle @ $B = (x_b, t_b)$
- From all A via Green's function (propagator) G

$$\psi(x_b,t_b) = \int dx_a G(x_b,t_b;x_a,t_a)\psi(x_a,t_a)$$

Huygen-Feynman Quantum Wavelets

Classical Becomes Quantum



- G(b; a) =spherical wavelet
- $\psi(x_b, t_b) = \sum$ wavelets

$$\psi(x_b, t_b) = \int dx_a G(b, a) \psi(x_a, t_a)$$

$$G(b,a) = \frac{\exp\left[i\frac{m(x_b - x_a)^2}{2(t_b - t_a)}\right]}{\sqrt{2\pi i(t_b - t_a)}}$$

- F's vision: $\psi \leftrightarrow \text{path}$
- $\psi_B = \sum_{\text{all}} \text{paths, A}$
- Δ paths Δ probabilities
- All paths possible!
- Also relativity, fields

Hamilton's Principle of Least Action (Classical)

Newton's Law $\equiv \delta S[\bar{x}(t)] = 0$

"The most general motion of a physical particle moving along the classical trajectory $\bar{x}(t)$ from time t_a to t_b is along a path such that the action $S[\bar{x}(t)]$ is an extremum."

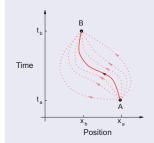
$$\delta S = S[\bar{x}(t) + \delta x(t)] - S[\bar{x}(t)] = 0 \quad (1)$$

(Constraint)
$$\delta(x_a) = \delta(x_b) = 0$$
 (2)

$$[x(t)] = functional$$

$$S[\bar{x}(t)] = \int_{t_{-}}^{t_{b}} dt \, L[x(t), \dot{x}(t)] \quad (3)$$

$$L = \text{Lagrangian} = T[x, \dot{x}] - V[x]$$
 (4)



Connecting CM Hamilton's Prin to QM Paths

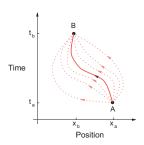
Consider Free Particle (V = 0)

$$S[b,a] = \int_{t_a}^{t_b} dt \, (T-V)$$

$$= \frac{m}{2}\dot{x}^2(t_b - t_a) = \frac{m}{2}\frac{(x_b - x_a)^2}{t_b - t_a}$$
 (1)

$$\Rightarrow G(b,a) = \frac{e^{iS[b,a]/\hbar}}{\sqrt{2\pi i(t_b - t_a)}}$$
 (2)

- F: QM = path integrals
- All paths ∃, △ prob
- Mainly classical



$$\Rightarrow G(b,a) = \sum_{\text{paths}} e^{iS[b,a]/\hbar} \quad (3)$$

•
$$\hbar \simeq 10^{-34} \, \mathrm{Js} \ \Rightarrow \ \sim \bar{x}$$

•
$$S_{\bar{x}}$$
 = extremum



Relate Paths to Ground State Wave Function

$\overline{\mathsf{Hermitian}\ ilde{H}} \Rightarrow \mathsf{Complete}\ \mathsf{Orthonormal}\ \mathsf{Set}\ \Rightarrow \mathsf{Propagator}$

$$\tilde{H}\psi_n = E_n\psi_n \tag{1}$$

$$\psi(x,t) = \sum_{n=0}^{\infty} c_n e^{-iE_n t} \psi_n(x)$$
 (2)

$$c_n = \int_{-\infty}^{+\infty} dx \, \psi_n^*(x,0) \psi(x,0) \tag{3}$$

$$\rightarrow \underline{\psi(x,t)} = \int_{-\infty}^{+\infty} dx_0 \sum_n \psi_n^*(x_0) \psi_n(x) e^{-iE_n t} \underline{\psi(x_0,t=0)}$$

Recall:
$$\psi(x_b, t_b) = \int dx_a G(x_b, t_b; x_a, t_a) \psi(x_a, t_a)$$
 (5)

$$\Rightarrow G(x, t; x_0, 0) = \sum_{n} \psi_n^*(x_0) \psi_n(x) e^{-iE_n t}$$
 (6)

Relate Space-Time Paths to Ψ_0 (cont)

Hermitian $\tilde{H} \Rightarrow$ Complete Orthonormal Set

$$G(x, t; x_0, t = 0) = \sum_{n} \psi_n^*(x_0) \psi_n(x) e^{-iE_n t}$$
 (1)

• Evaluate @ imaginary t (Wick rotation):

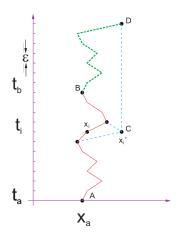
$$G(x, -i\tau; x_0, t = 0) = \sum_{n} \psi_n^*(x_0) \psi_n(x) e^{-E_n \tau}$$
 (2)

- Im time $\tau \to \infty$ only n = 0
- For $|\psi_0|^2$: paths start & end at $x_0 = x$

$$G(x,-i\tau;x,0) = \sum_{n} |\psi_n(x)|^2 e^{-E_n\tau} = |\psi_0|^2 e^{-E_0\tau} + |\psi_1|^2 e^{-E_1\tau} + \cdots$$
 (3)

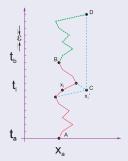
$$\Rightarrow |\psi_0(x)|^2 = \lim_{\tau \to \infty} e^{E_0 \tau} G(x, -i\tau; x, 0)$$
 (4)

Break Now, Compute Later



Lattice Quantum Mechanics (Algorithm)

Easy: Discrete Times & Positions Only!



- Path: ∑ links
- Euler + Time step ε :

$$\frac{dx_j}{dt} \simeq \frac{x_j - x_{j-1}}{\varepsilon} \tag{1}$$

$$S_j \simeq L_j \Delta t$$
 (2)

$$\simeq \frac{m\Delta x^2}{2\varepsilon} - V(x_j)\varepsilon$$
 (3)

- Add actions for N-links
- $G(b, a) \leftrightarrow \sum_{a-b \text{ paths}}$
- Ea path = \sum_{links}

$$G(b,a) = \int dx_1 \cdots dx_{N-1} e^{iS[b,a]}$$
 (4)

Rotate t: Lagrangian $(-i\tau) = -Hamiltonian (\tau)$

Wick Rotation into Imaginary Time

$$G(x, t; x_0, t_0) = \sum dx_1 dx_2 \cdots dx_{N-1} e^{iS[x, x_0]}$$
 (1)

$$S[x, x_0] \simeq \sum_{j=1}^{N-1} L(x_j, \dot{x}_j) \varepsilon$$
 (2)

$$L(x,\dot{x}) = T - V(x) = +\frac{1}{2}m\left(\frac{dx}{dt}\right)^2 - V(x)$$
 (3)

$$\Rightarrow L\left(x, \frac{dx}{-id\tau}\right) = -\frac{1}{2}m\left(\frac{dx}{d\tau}\right)^2 - V(x) = -H \tag{4}$$

Put Pieces Together Sum Over Paths

Related to Wave Function

$$G(x, -i\tau; x_0, 0) = \int dx_1 \dots dx_{N-1} e^{-\int_0^{\tau} H(\tau')d\tau'}$$
 (5)

Individual path integral:

$$\int H(\tau) d\tau \simeq \sum_{j} \varepsilon E_{j} = \varepsilon \mathcal{E}$$
 (6)

Ground State Wave Function via Feynman:

$$|\psi_0(x)|^2 = \frac{1}{Z} \lim_{\tau \to \infty} \int_{\text{naths}} dx_1 \cdots dx_{N-1} e^{-\varepsilon \mathcal{E}}$$
 (7)

QM Paths ψ_0 Lattice t
ightarrow -i au Trick Implementation Assessment

Imaginary Time Relates QM to Thermodynamics

Schrödinger Equation → Heat Diffusion Equation

• t in QM $\rightarrow -i\tau$

$$i\frac{\partial \psi}{\partial (-i\tau)} = \frac{-\nabla^2}{2m}\psi \quad \Rightarrow \quad \frac{\partial \psi}{\partial \tau} = \frac{\nabla^2}{2m}\psi \tag{1}$$

- Boltzmann $\mathcal{P} = e^{-\varepsilon \mathcal{E}}$ weights ea Feynman path
- Temperature ⇔ time step:

$$\mathcal{P} = \mathbf{e}^{-\varepsilon \mathcal{E}} = \mathbf{e}^{-\varepsilon/k_B T} \quad \Rightarrow \quad k_B T = \frac{1}{\varepsilon} \equiv \frac{\hbar}{\varepsilon} \tag{2}$$

- $\bullet \Rightarrow \lim_{\varepsilon \to 0} = \lim_{T \to \infty}$
- ψ_0 : long imaginary τ vs $\hbar/\Delta E$
- Like equilibration in Ising model

Summary (This is Heavy Stuff)

Feynman's Path Integral Formulation of QM

- QM ψ via statistical fluctuations \sim class trajectory
- Propagator($t_a \rightarrow t_b$) G = path integral, $\sum_{\text{paths}} \int$
- Hamilton: extremum $S \rightarrow$ path integration of H
- Path integral = sum trajectories on x-t lattice
- Paths weighted with probability $e^{-iS/\hbar}$
- Algorithm: \triangle path link $\Rightarrow \triangle E$ (like Ising)
- Ψ equilibrates to ground state

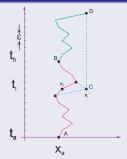
Break Before Algorithm

Quantum Monte Carlo (QMC) Applet



A Time-Saving Trick

Compute $\psi(x)$ for All $x(x_b)$ Simultaneously



- Integrate all x sites
- Don't compute $\delta(x)$!
- Accumulate $\psi(x')$

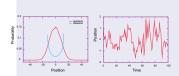
$$\begin{aligned} |\psi_0(x)|^2 &= \int dx_1 \cdots dx_N \, e^{-\varepsilon \mathcal{E}(x,x_1,\ldots)} \\ &= \int dx_0 \cdots dx_N \delta(x-x_0) \, e^{-\varepsilon \mathcal{E}(\ldots)} \end{aligned}$$

- Frequent $x_i \Rightarrow \text{larger } \psi(x_i)$
- EG: AB, New path + C
- *CBD* same $\sum E_i$ as *ACB*
- Equilibrate, flip links, new E

Lattice Implementation

QMC.py





Harmonic oscillator

$$V(x) = \frac{1}{2}x^2$$

- Natural units: m = 1, L: $\sqrt{\hbar/m\omega}$; t: $1/\omega$; $T = 2\pi$
- 3 Short $T \sim 2T$, Long $t \sim 20T$
- **1** Classical: max ρ @ turning pts
- **5** Each x_i , running sum $|\Psi_0(x_i)|^2$
- \bullet Δ seed; many runs > 1 long run

Assessment and Exploration

- Plot classical trajectory, some actual space-time paths
- 2 Explore effect of smaller Δx , smaller Δt
- **3** Assume $\psi(x) = \sqrt{\psi^2(x)}$, calculate:

$$E = \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle} = \frac{\omega}{2 \langle \psi | \psi \rangle} \int_{-\infty}^{+\infty} \psi^*(x) \left(\frac{-d^2}{dx^2} + x^2 \right) \psi(x) dx \qquad (1)$$

- Explore effect of larger, smaller ħ
- **1** Test ψ with quantum bouncer:

$$V(x) = mg|x| \tag{2}$$

$$x(t) = x_0 + v_0 t + \frac{1}{2}gt^2.$$
(3)

QM Paths ψ_0 Lattice t
ightarrow -i au Trick Implementation Assessment

Summary

Feynman Path Integrals



- A different view of quantum mechanics
- It seems to give same answers as traditional QM
- Is at heart of lattice quantum chromodynamics
- Hard to apply beyond ground state
- Satisfying connection to classical mechanics