

Integral Equations in Quantum Mechanics II

I Bound States, II Scattering*

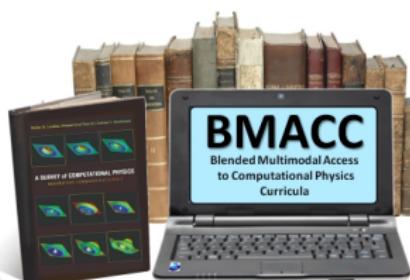
Rubin H Landau

Sally Haerer, Producer-Director

Based on *A Survey of Computational Physics* by Landau, Páez, & Bordeianu

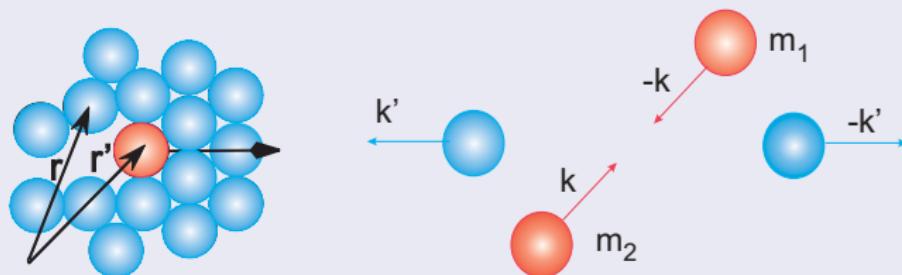
with Support from the National Science Foundation

Course: **Computational Physics II**



Problem: Nonlocal Potential Scattering

Integro-Differential Equation



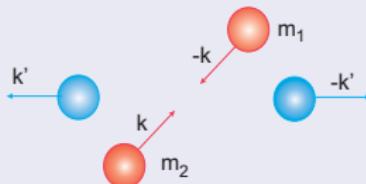
- **Problem:** Scattering from nonlocal potential

$$-\frac{1}{2m} \frac{d^2\psi(r)}{dr^2} + \int dr' V(r, r')\psi(r') = E\psi(r) \quad (1)$$

- Avoid Integro-Differential Equation

Theory: Lippmann–Schwinger Equation

Integral Form Schrödinger Equation



- Better solve scattering amplitudes = observable ($\neq \psi$)
- \mathcal{P} = Cauchy principal-value prescription; $l = 0$, $\hbar = 1$

$$R(k', k) = V(k', k) + \frac{2}{\pi} \mathcal{P} \int_0^\infty dp \frac{p^2 V(k', p) R(p, k)}{(k_0^2 - p^2)/2\mu} \quad (1)$$

$$E = \frac{k_0^2}{2\mu}, \quad \mu = \frac{m_1 m_2}{m_1 + m_2} \quad (2)$$

Math: Computing Singular Integrals

How Handle Singularity?

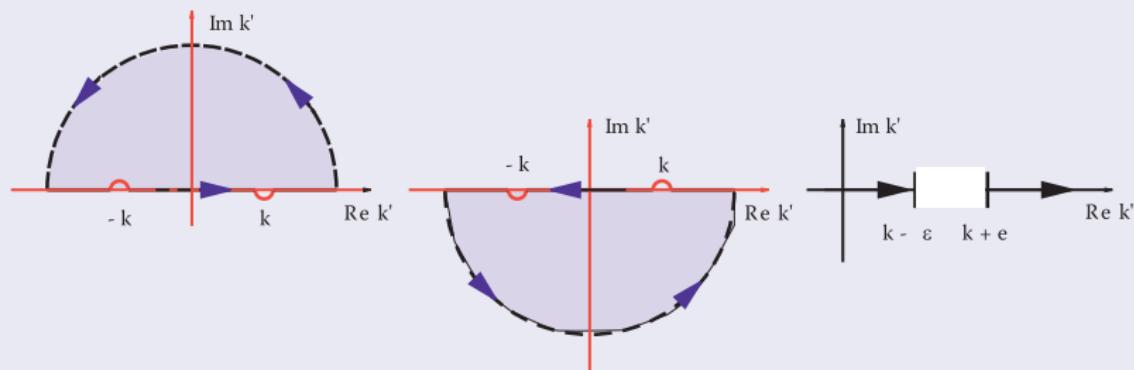
- **Singular** integral \mathcal{G} has singular integrand $g(k)$:

$$\mathcal{G} = \int_a^b g(k) dk \neq \infty \quad (1)$$

- Computational **danger** near singularity

Math: Computing Singular Integrals

What to Do at Singularity?



- Give energy k_0 small imaginary part $\pm i\epsilon$
- Cauchy principal-value $\mathcal{P} = \text{pinch}^*$

$$\mathcal{P} \int_{-\infty}^{+\infty} f(k) dk = \lim_{\epsilon \rightarrow 0} \left[\int_{-\infty}^{k_0 - \epsilon} f(k) dk + \int_{k_0 + \epsilon}^{+\infty} f(k) dk \right] \quad (1)$$

Numerical Principal Values

Subtract Zero Integral (**Hilbert Transform**)

$$\mathcal{P} \int_{-\infty}^{+\infty} \frac{dk}{k - k_0} = \mathcal{P} \int_0^{+\infty} \frac{dk}{k^2 - k_0^2} = 0 \quad (1)$$

$$\mathcal{P} \int_0^{+\infty} \frac{f(k) dk}{k^2 - k_0^2} = \int_0^{+\infty} \frac{[f(k) - f(k_0)] dk}{k^2 - k_0^2} \quad (2)$$

- NB No RHS \mathcal{P} , no $k = k_0$ singularity

Method: Integral- to Linear- to Matrix-Equations

Rewrite Principal-Value as Definite Integral

$$R(k', k) = V(k', k) + \frac{2}{\pi} \int_0^\infty dp \frac{p^2 V(k', p) R(p, k) - k_0^2 V(k', k_0) R(k_0, k)}{(k_0^2 - p^2)/2\mu} \quad (1)$$

- Convert to linear equations; approximate integral:

$$\begin{aligned} R(k, k_0) &\simeq V(k, k_0) + \frac{2}{\pi} \sum_{j=1}^N \frac{k_j^2 V(k, k_j) R(k_j, k_0) w_j}{(k_0^2 - k_j^2)/2\mu} \\ &\quad - \frac{2}{\pi} k_0^2 V(k, k_0) R(k_0, k_0) \sum_{m=1}^N \frac{w_m}{(k_0^2 - k_m^2)/2\mu} \end{aligned} \quad (2)$$

- $(N+1)$ unknown $R(k_j, k_0), j = 0, N$

Method: Integral- to Linear- to Matrix-Equations

Evaluate for $k = N$ Gauss Points $k +$ Experimental k_0

$$k = k_i = \begin{cases} k_j, & j = 1, N \quad (\text{quadrature points}), \\ k_0, & i = 0 \quad (\text{experimental point}) \end{cases} \quad (1)$$

$$R_i = V_i + \frac{2}{\pi} \sum_{j=1}^N \frac{k_j^2 V_{ij} R_j w_j}{(k_0^2 - k_j^2)/2\mu} - \frac{2}{\pi} k_0^2 V_{i0} R_0 \sum_{m=1}^N \frac{w_m}{(k_0^2 - k_m^2)/2\mu} \quad (2)$$

- Express as matrix equations:

$$D_i = \begin{cases} +\frac{2}{\pi} \frac{w_i k_i^2}{(k_0^2 - k_i^2)/2\mu}, & \text{for } i = 1, N, \\ -\frac{2}{\pi} \sum_{j=1}^N \frac{w_j k_0^2}{(k_0^2 - k_j^2)/2\mu}, & \text{for } i = 0 \end{cases} \quad (3)$$

$$R = V + DVR \Rightarrow R = (1 - DV)^{-1} V \quad (4)$$

Solution via Matrix Inversion, Gaussian Elimination



CODE

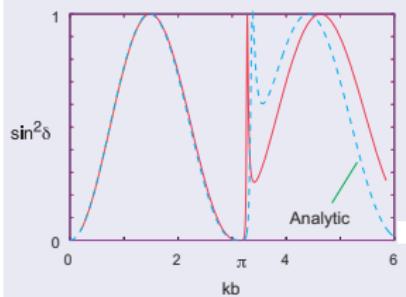
$$R = V + DVR \quad (1)$$

$$R = (1 - DV)^{-1} V \quad (2)$$

- Matrix inversion = direct, not fastest
- Matrix inversion = standard in mathematical libraries
- Useful if need $[1 - DV]^{-1}$
- Else **Gaussian elimination**

Implementation: Delta-Shell Potential

$\sin^2 \delta_0 \propto l = 0$ Cross Section



$$V(k', k) = \frac{-|\lambda| \sin(k'b) \sin(kb)}{2\mu k' k} \quad (1)$$

- Check analytic phase shift:

$$\tan \delta_0 = \frac{\lambda b \sin^2(kb)}{kb - \lambda b \sin(kb) \cos(kb)} \quad (2)$$

$$R(k_0, k_0) = -\frac{\tan \delta}{2\mu k_0} \quad (3)$$

- Estimate precision by increasing N grid points ($N = 26$)